

The Einstein Coefficients



Let N_1 and N_2 be the number of atoms per unit volume present in the energy levels E_1 and E_2 , respectively. Let $u(\omega)d\omega$ be the energy per unit volume of the incident radiation between frequency ω and $\omega + d\omega$.

$$\omega = (E_2 - E_1) / \hbar$$

(i) An atom in the lower energy level E_1 can absorb the incident radiation and go to the higher energy level E_2 . The number of atoms undergoing absorptions per unit time per unit volume is given by

$$R_{abs} = B_{12}N_1u(\omega) \quad (1)$$

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(ii) In a reverse process, an atom in the higher energy level E_2 can emit radiation at frequency ω and go to the level E_1 . This is stimulated emission. The number of atoms undergoing stimulated emissions per unit time per unit volume is given by

$$R_{st} = B_{21}N_2u(\omega) \quad (2)$$

(iii) An atom in the upper state E_2 can also make a spontaneous emission at frequency ω and go to the level E_1 . The number of atoms undergoing spontaneous emissions per unit time per unit volume is given by

$$R_{sp} = AN_2 \quad (3)$$

At thermal equilibrium, the number of upward transitions should be equal to downward transitions. If they are not equal, radiation energy will grow or sink which is not consistent with the assumption of thermal equilibrium.

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$$AN_2 + B_{21}N_2u(\omega) = B_{12}N_1u(\omega)$$

$$u(\omega) = \frac{A}{(N_1/N_2)B_{12} - B_{21}} \quad (4)$$

Using Boltzmann's law, under equilibrium, the ratio of N_2 and N_1 is

$$\frac{N_2}{N_1} = e^{-(E_2-E_1)/k_B T} = e^{-\hbar\omega/k_B T} \quad (5)$$

k_B = Boltzmann's constant = 1.38×10^{-23} J/K.

According to Planck's theory, the radiation energy density per unit frequency interval is

$$u(\omega) = \frac{\hbar\omega^3 n_0^3}{\pi^2 c^3} \frac{1}{e^{\hbar\omega/k_B T} - 1} \quad (6)$$

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Comparing Eq. (4) and (6), we see that

$$B_{21} = B_{12} \equiv B$$

$$A_{21} = \frac{B\hbar\omega^3 n_0^3}{\pi^2 c^3}$$

The coefficients A and B are called Einstein coefficients.

Observations :

(i) The stimulated emission rate per atom is same as the absorption rate per atom.

(ii) At thermal equilibrium, the ratio of spontaneous to stimulated emission is

$$R = \frac{AN_2}{BN_2u(\omega)} = e^{\hbar\omega/k_B T} - 1$$

If $\hbar\omega \gg k_B T$, $R \gg 1$ and spontaneous emission dominates the stimulated emission.

Example 1

Consider an optical source at $T = 1000\text{K}$. Calculate the ratio of the number of spontaneous to stimulated emission for a wavelength of

1.55 microns. Assume $k_B = 1.38 \times 10^{-23} \text{J/K}$, $\hbar = 1.054 \times 10^{-34} \text{Js}$.

$\lambda = 1.55 \text{ microns}$ implies $f = c / \lambda = 3 \times 10^8 / (1.55 \times 10^{-6}) = 1.9355 \times 10^{14} \text{ Hz}$

The ratio of the number of spontaneous to stimulated emission is given by

$$R = \exp(\hbar 2\pi f / k_B T) - 1$$

$$= \exp(1.054 \times 10^{-34} \times 2 \times \pi \times 1.9355 \times 10^{14} / (1.38 \times 10^{-23} \times 10^3)) - 1$$

$$= 1.08 \times 10^4$$

Therefore at optical frequencies and at this temperature, spontaneous emission dominates and hence the light from usual light sources is incoherent

Population inversion

Light amplification can take place only if $R_{st} > R_{abs}$ or

$$N_2 > N_1.$$

Under thermal equilibrium, it is $N_1 > N_2$. Therefore, N_2 should be increased by external means. The condition $N_2 > N_1$ is referred as *population inversion*.